



Cherenkov Radiation and Its Application

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Brief history of Cherenkov rad.

- Discovered by Pavel Cherenkov 1934
- when studying fluorescence emission induced in uranyl salts
(γ -ray in radioactive source)
- light was emitted even with solvent only
- Explained Ilya Frank & Igor Tamm 1937

Cherenkov radiation theory

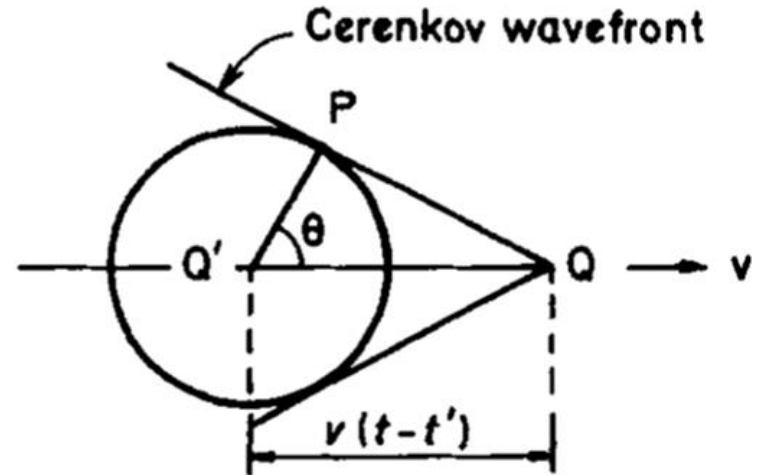
What is Cherenkov Radiation (CR)

- emitted by a high energy charged particle
- with a velocity greater than the light speed in medium
- usually form a forward light cone
- occur when the condition is met:

$\beta n > 1$

v/c , ratio of particle velocity to the vacuum light speed

refractive index of medium



A geometric proof of CR

(particle travel from t' to t)

- light path $Q'P = \frac{c}{n}(t - t')$
- Particle path $Q'Q = v(t - t')$,
- Emission angle $\cos \theta = \frac{Q'P}{Q'Q} = \frac{c}{vn} = \frac{1}{\beta n}$
- $|\cos \theta| \leq 1 \rightarrow \beta n > 1$

travel from t' to t)

path $Q'P = \frac{c}{n}(t - t')$
 particle path $Q'Q = v(t - t')$,
 emission angle $\cos \theta = \frac{Q'P}{Q'Q} = \frac{c}{vn} = \frac{1}{\beta n}$
 $|\cos \theta| \leq 1 \rightarrow \beta n > 1$

CR field in EM theory

- In vacuum: E-field at a field point (time t , charge q , uniform velocity $\mathbf{v}$)

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \frac{\alpha \mathbf{R}}{s^3},$$

where $\mathbf{R}$: radius vector from charge location at time t to the field point,

$$\alpha = \left(1 - \frac{v}{c}\right)^2, \quad s = \left[\alpha R^2 + \frac{1}{c^2} (\mathbf{v} \cdot \mathbf{R})^2\right]^{\frac{1}{2}}$$

- In medium (permittivity ϵ & refractive index n),

$$\alpha \Rightarrow \left(1 - \frac{vn}{c}\right)^2, \quad s \Rightarrow \left[\alpha R^2 + \left(\frac{n}{c} \mathbf{v} \cdot \mathbf{R}\right)^2\right]^{\frac{1}{2}}$$

CR field in EM theory

$$s = \left[1 - \left(\frac{vn}{c} \right)^2 \sin^2 \varphi \right]^{\frac{1}{2}} R, \text{ where } \varphi \text{ is the angle between } \mathbf{v} \text{ and } \mathbf{R}$$

- If $v > \frac{c}{n}$, $s \Rightarrow$ imaginary, except for region with $\sin \varphi \leq \frac{c}{vn}$
- Field point must be to the rear of particle
- E-field exist only in a cone, with half angle $\varphi = \arcsin\left(\frac{c}{vn}\right)$, and vertex at the particle location at time t .
- Cone surface = surface of Cherenkov shock wave and contains the Cherenkov radiation field. ($E \rightarrow \infty$ on the surface)

CR field in EM theory



- A unit vector $\mathbf{S}$: normal to the upper plane of light cone, and
- $\mathbf{S} \parallel \mathbf{k}$ ($|\mathbf{k}| = \frac{\omega}{c}n$) of the Cherenkov radiation. Then

$$k_x = \frac{\omega}{c}n \sin \theta = \frac{\omega}{c}n \left(\frac{c}{nv} \right) = \frac{\omega}{v}, \text{ and}$$

$$k_{yz} = \sqrt{\left(\frac{\omega}{c}n \right)^2 - \left(\frac{\omega}{v} \right)^2}$$

- If $v < \frac{c}{n}$, k_{yz} will become imaginary $\Rightarrow$ The field will not exist

Threshold particle energy

- $E(\text{particle}) = \gamma m_0 c^2$

$$= \frac{1}{\sqrt{1-\beta^2}} m_0 c^2$$

$$\geq \frac{n}{\sqrt{n^2-1}} m_0 c^2$$

- For example, in water, $E_{\text{thr}} \sim 264 \text{ keV}$

in air , $E_{\text{thr}} \sim 38 \text{ GeV}$

Cherenkov counter

Cherenkov counter

Made by (fitting a long pipe of) medium with an optical system capable of

- detecting the emitted light
- measuring the angle of emission.

Usage:

- With known rest mass/type,
- Measurement of the Cherenkov angle $\Rightarrow$ Measurement of the particle momentum/speed.
- Or vice versa

Imaging atmospheric Cherenkov telescope (IACT)

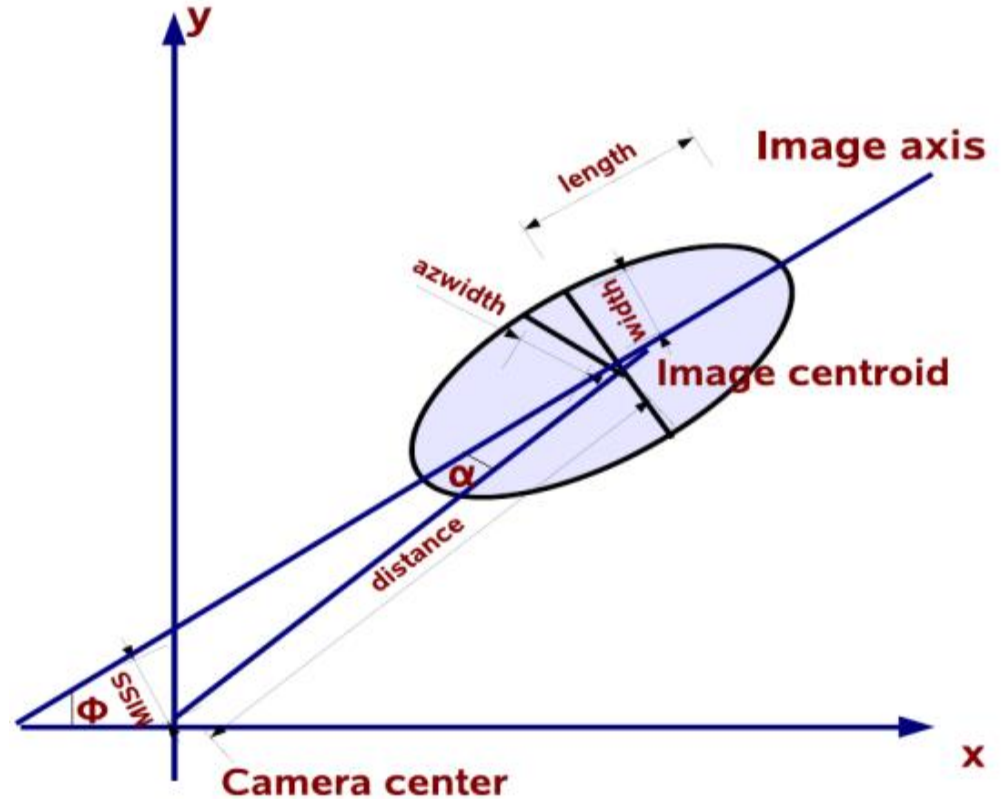


Working principles

- aims at γ -ray with 100GeV-20TeV
- use the atmosphere as the medium
- high-energy γ -ray interact and scatter $\rightarrow$ produce a shower of charged particles with $v > c_{\text{air}}$
- Shower produces Cherenkov pool with $r > 100\text{m}$
- Mirrors on a telescope reflect CR towards a fast-camera (ns) and images are obtained
- light will be captured by sensors and turned into data for later analysis.

Hillas parameters

- Size: sum of photo-electrons (pe) of all PMT channel
- Distance: angular distance of centroid from telescope axis
- Alpha: angle between major axis and telescope axis
- useful in simulations, γ /hadrons segregation



Computer simulation

tool: CORSIKA code

- tracks each particale produced in shower, not only γ -ray
- outputs particle type, energy, direction etc which can be used in further analysis and simulations

with add-on package,

- able to extract cherenkov photon data from shower data

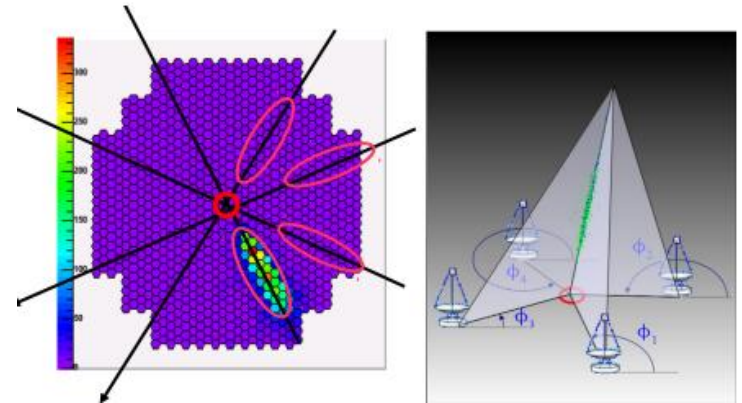
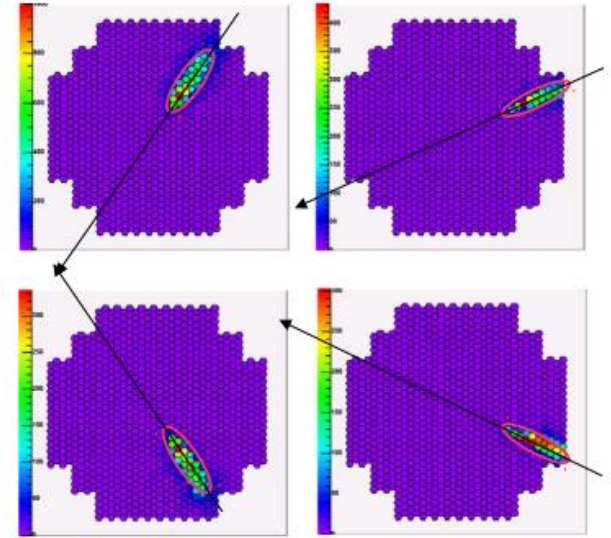
Primary γ -ray direction estimation

For multiple telescopes

- major axis is direction of primary γ -ray projecting onto camera plane
- so γ -ray source lies on major axis
- finding the intersection of major axes

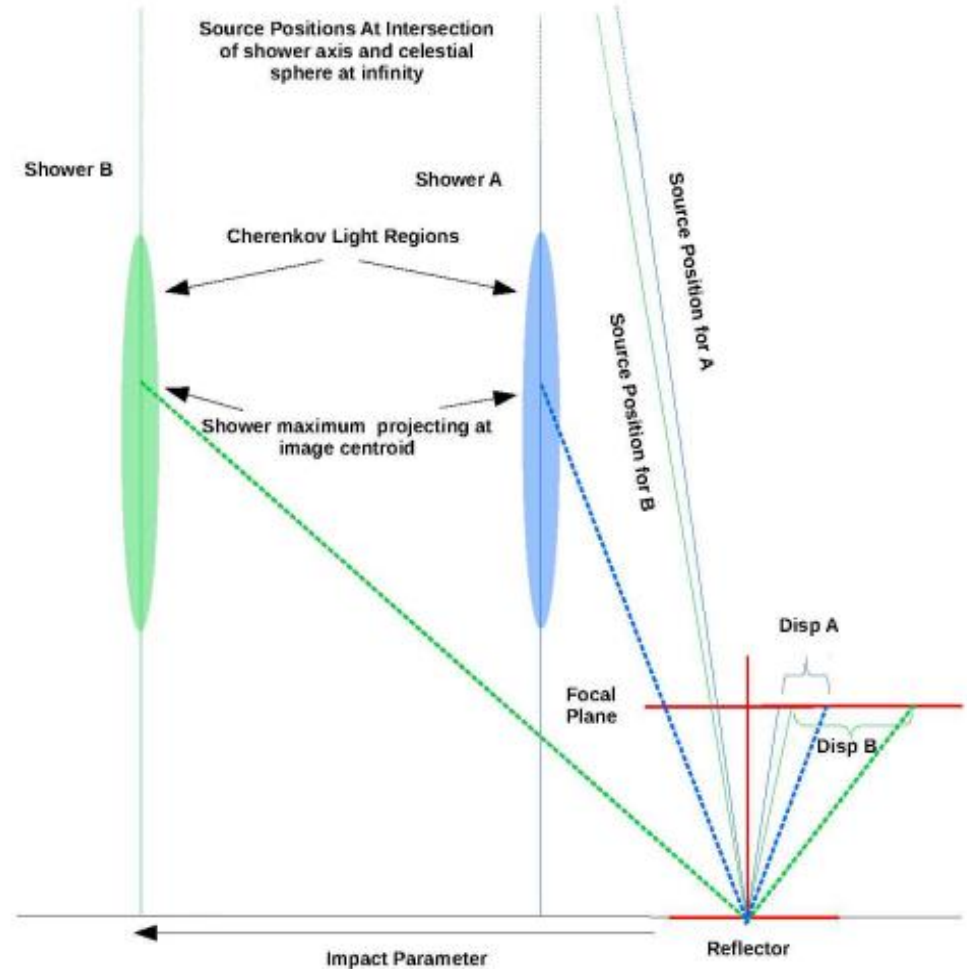
gives the source position on camera

For single telescope, Disp method is a way



Disp method

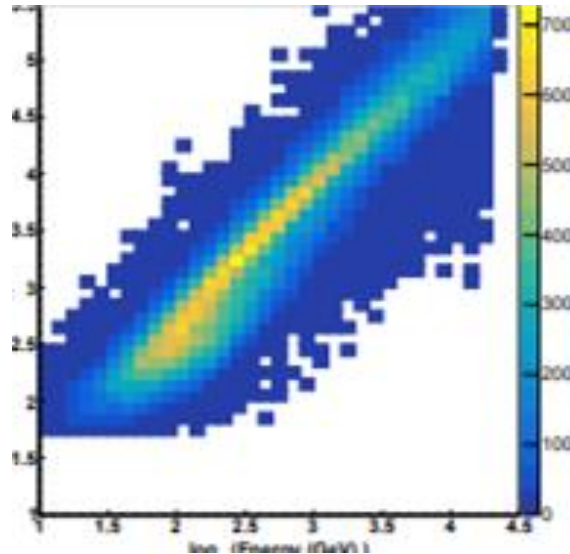
- angular distance between source position and image centroid
- disp increases with impact parameter, same for elongation (width / length)
- disp-elongation correlation can be estimated by a computer model (RFM regression)



γ -ray energy estimation

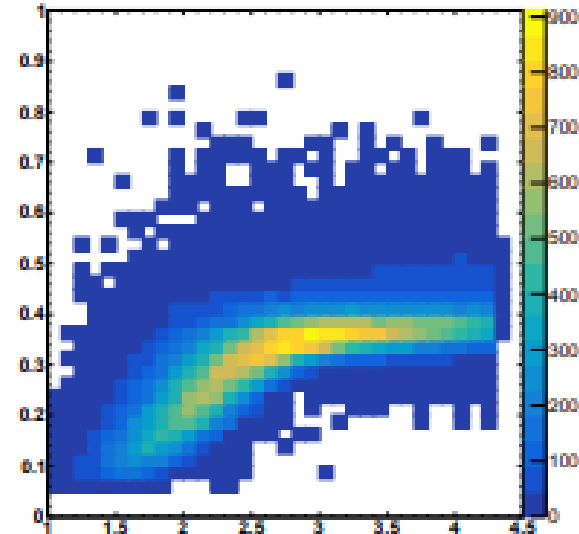
- primary energy relates to Hillar parameters (size, length) => allow estimation
- can also be estimated by RFM regression

log(size[pe])



log(energy[GeV])

length[degree]



log(energy[GeV])

Night sky light contribution (LONS)

- raw image of shower contains pe induced by LONS photons which affects sensitivity of telescope
- Measure LONS spectrum at nearby site
- To find the contribution of LONS on each pixel per unit time,

$$dN/dt (1 \text{ pixel}) = (\int dN/dtdA d\Omega/d\lambda d\lambda) \times A\Omega(1 \text{ pixel})$$

number of pe (per
pixel, LONS)

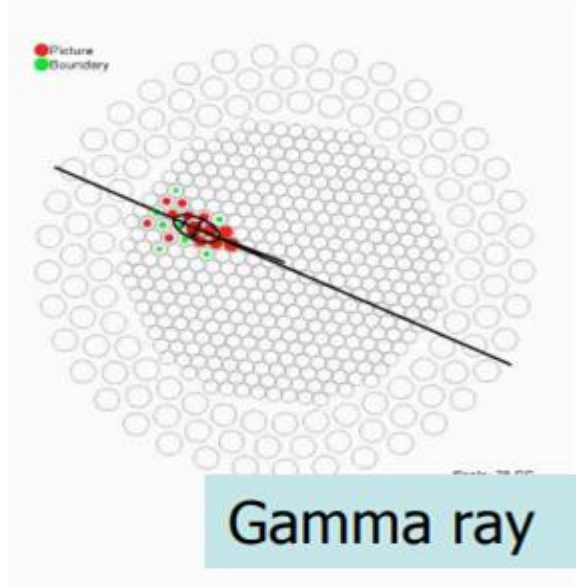
(1.46 pe per pixel in 5ns)

LONS flux spectrum

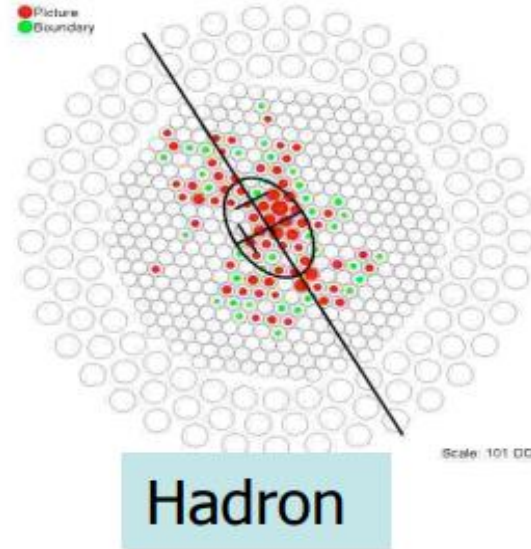
total reflector area

angular size of 1 pixel

Rejection of background cosmic ray



Shorter in width & length



longer in width & length

Rejection of background cosmic ray

Method 1: static cut

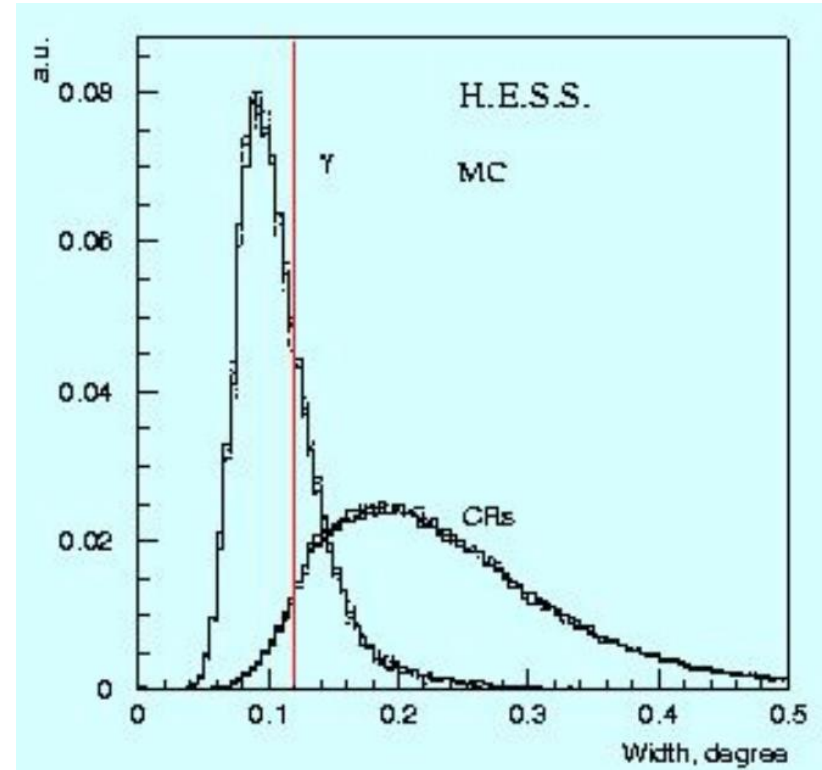
- apply a simple cut on Hillas parameters

More efficient way: Machine learning

- Trained using data from simulations
- Hillas parameters as input
- Assign a hadronness score

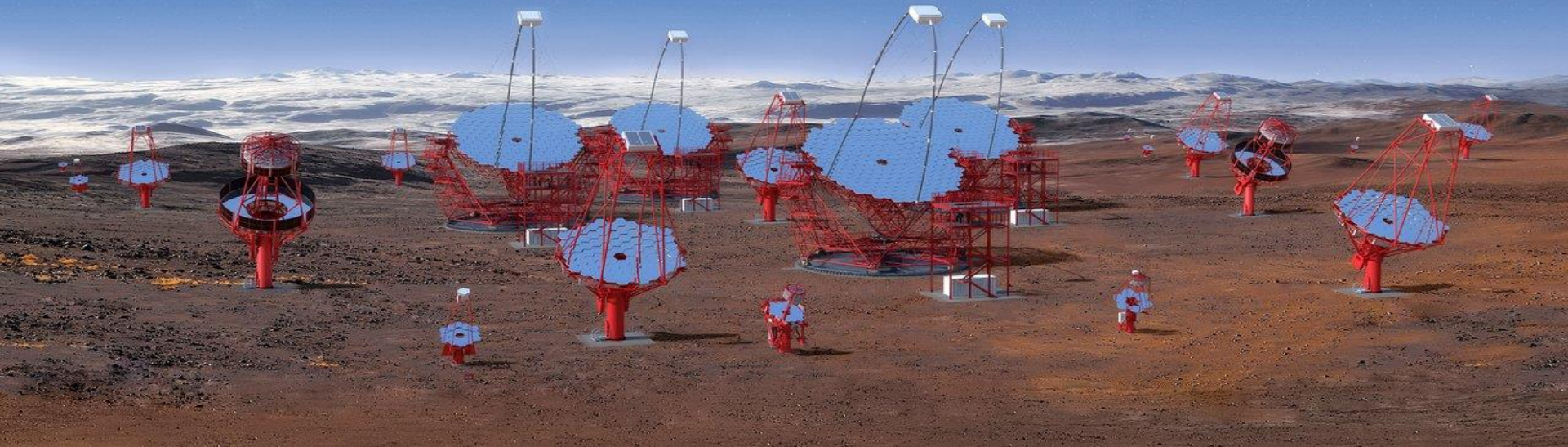
(likeliness of being a hadron event)

- A cut is applied to the hadronness



Static cut example

New generation γ detection telescope



Locations of CTA

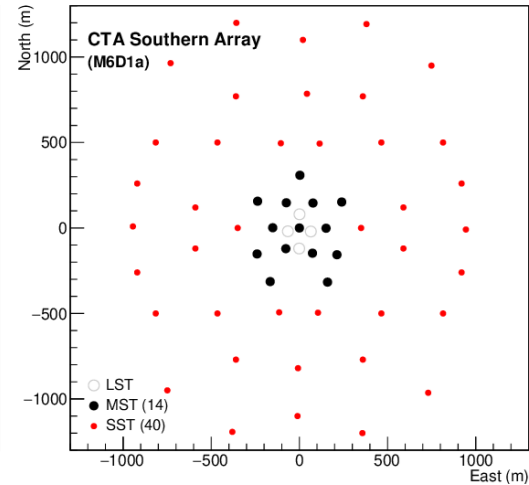
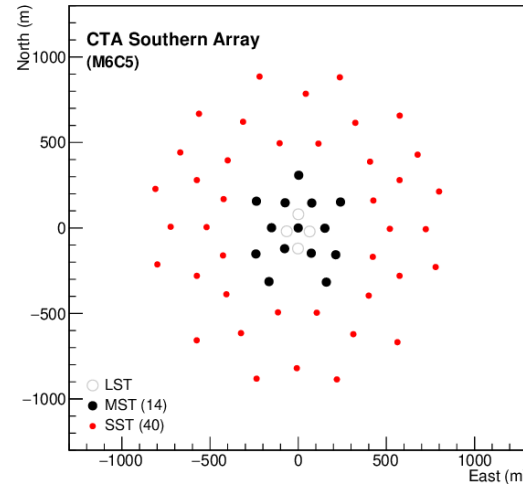
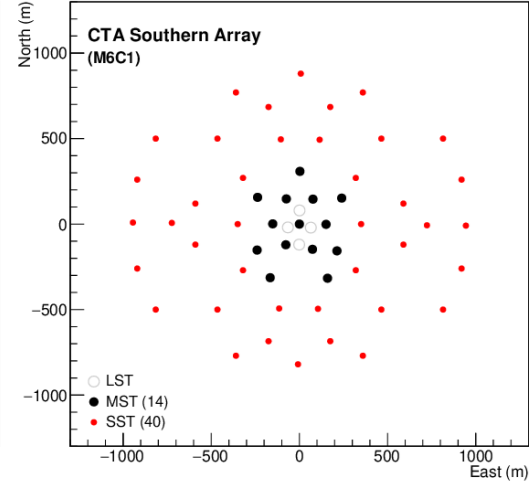
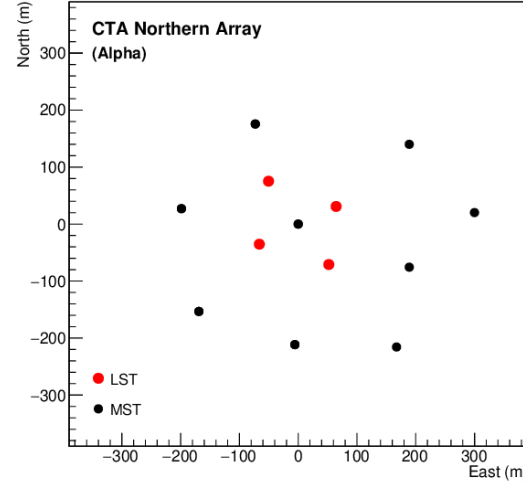


Types of telescopes

Large size telescopes: 20-100GeV

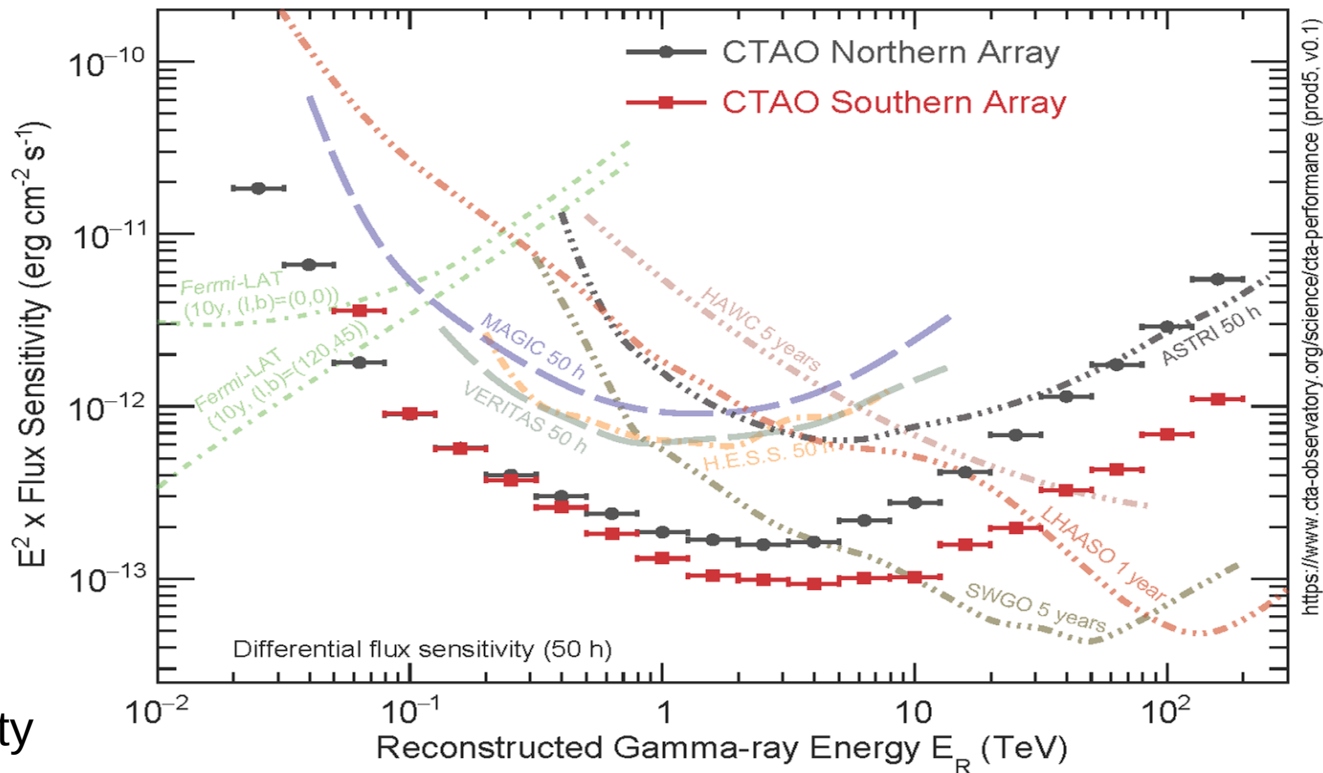
Medium size telescopes: 0.1-5TeV

Small size telescopes: 5-100TeV



CTA array arrangement

Comparison with CTA and other tools



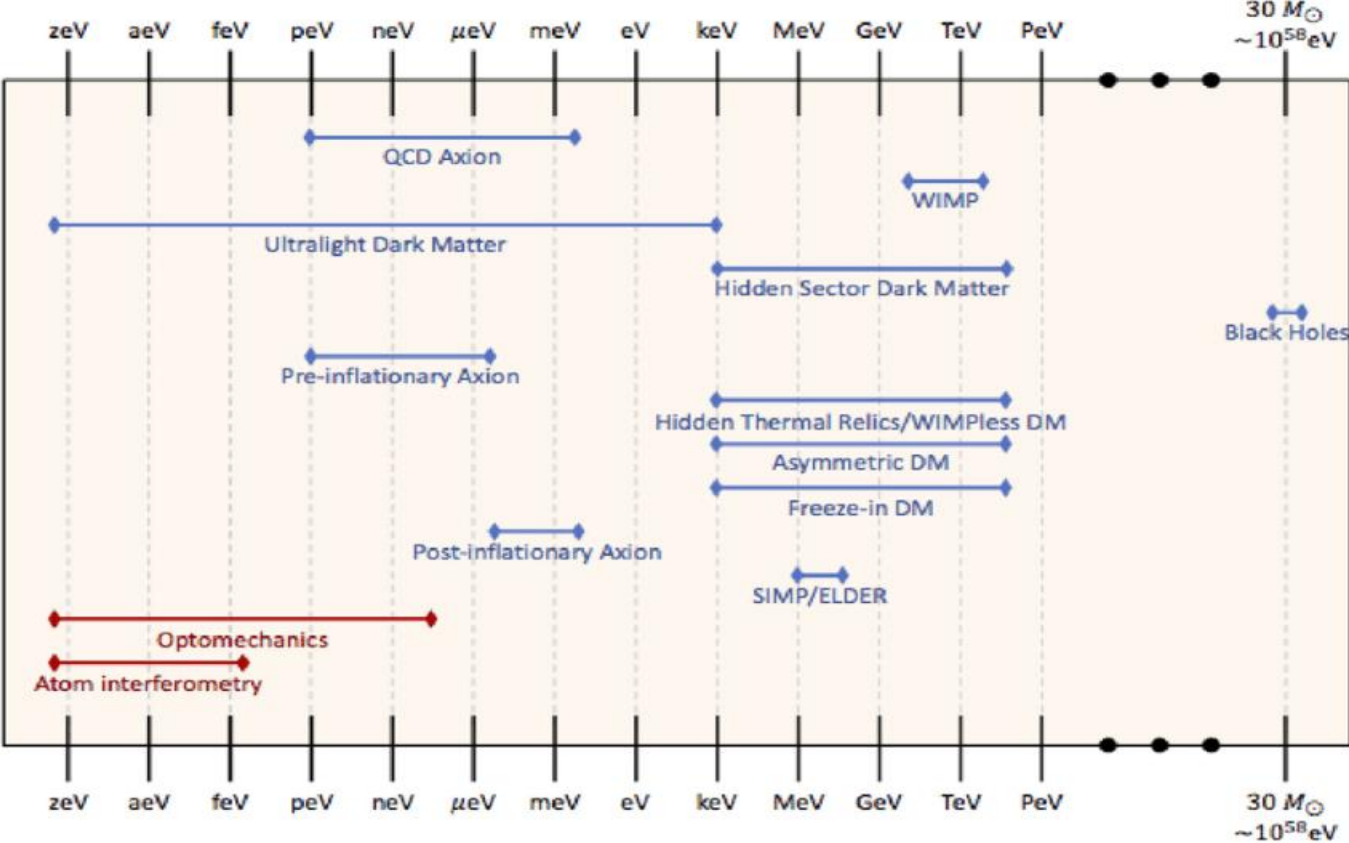
<https://www.cta-observatory.org/science/cta-performance> (prod5, v0.1)

- Sensitivity
- detection range

Dark matter (DM) detection with CR

Dark matter candidates

[\(PDF\) Quantum physics in space | ResearchGate](#)



Why WIMP?

$$\Omega_{\chi} h^2 = \frac{3 \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma_{\text{an}} v \rangle}$$

Cosmological abundance of WIMP

Hubble ratio = $H_0/100 \text{ km s}^{-1} \text{ Mpc}^{-1}$

Weak interaction cross-section, $3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

- $\Omega_{\chi} \sim 0.2$, similar to $\Omega_{\text{DM}} \sim 0.27$
- Main focus - Neutralino (self-annihilated)

Dark matter annihilation

$$\frac{d\Phi_\gamma}{d\Omega dE_\gamma}(E_\gamma, \psi) = \frac{1}{4\pi} \int_{\text{l.o.s}} d\ell(\psi) \rho_\chi^2(\mathbf{r}) \left(\frac{\langle \sigma v \rangle_{\text{ann}}}{2S_\chi m_\chi^2} \sum_f B_f \frac{dN_\gamma^f}{dE_\gamma} \right)$$

differential flux astrophysic factor particle physics factor

where,

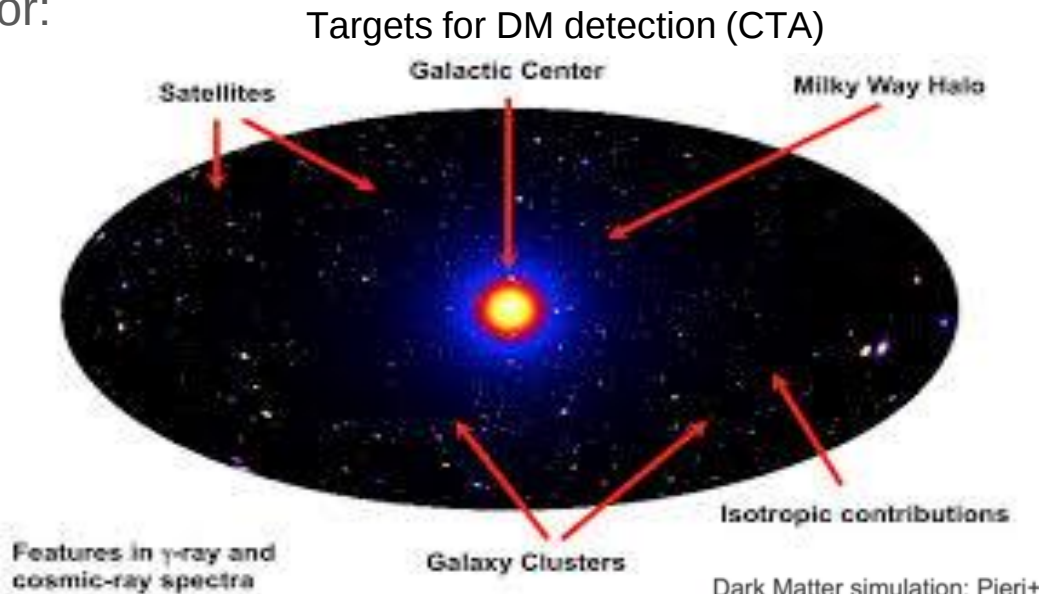
- σ = total cross section
- dN/dE is spectrum
- ρ is DM density
- $S_\chi = 1$ (self-annihilated) or 2
(else)

Good detection target

$$\frac{d\Phi_\gamma}{d\Omega dE_\gamma}(E_\gamma, \psi) = \frac{1}{4\pi} \int_{\text{l.o.s}} d\ell(\psi) \rho_\chi^2(\mathbf{r}) \left(\frac{\langle\sigma v\rangle_{\text{ann}}}{2S_\chi m_\chi^2} \sum_f B_f \frac{dN_\gamma^f}{dE_\gamma} \right)$$

To maximize γ flux, we need to look for:

- a close ($\propto d\Omega$),
 - DM-dominated system ($\propto \rho^2$)
- => Galactic centre



Dark Matter simulation: Pieri+
[2011PhRvD..83b3518P](#)

Noise suppression

Major noise: Strong γ -ray background (almost isotropic)

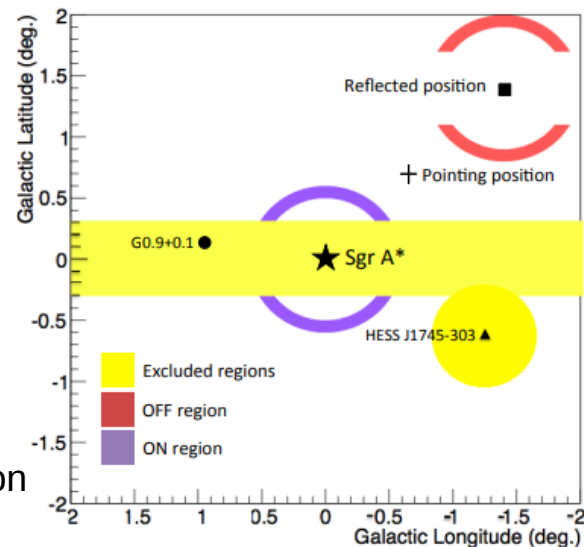
Solutions:

- Observe γ -ray background from other position that are believed without DM
- Build emission model to predict γ -ray emission without DM

Also, when detecting galactic center,

- exclude areas around the galactic core, as γ -flux is too high.

(by massive stars, neutron stars etc.)



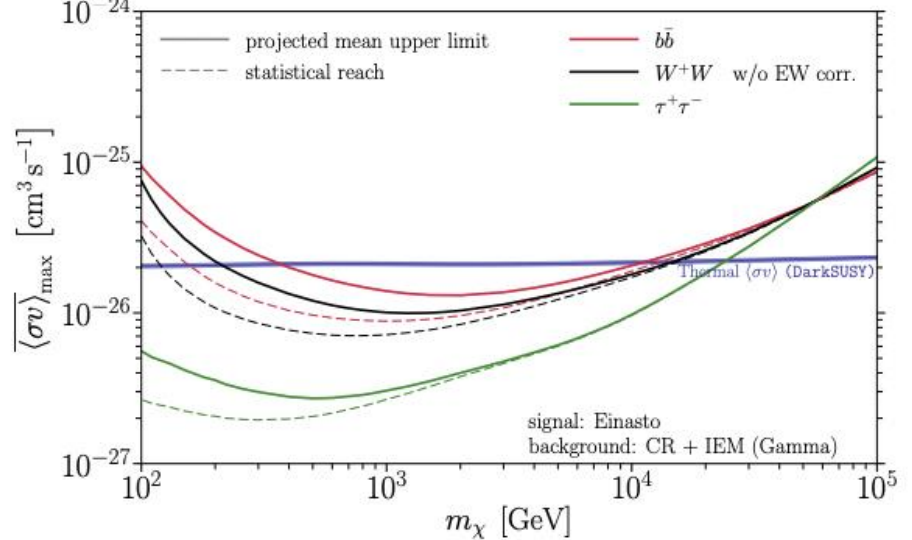
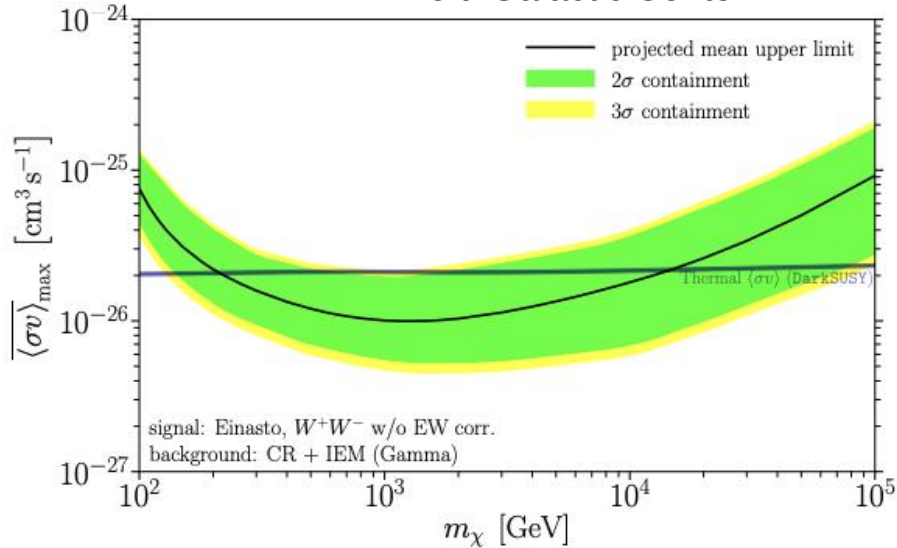
Galactic center detection
region (HESS)

Sensitivity of CTA in DM detection

Sensitivity curves to a DM annihilation signal for different DM model

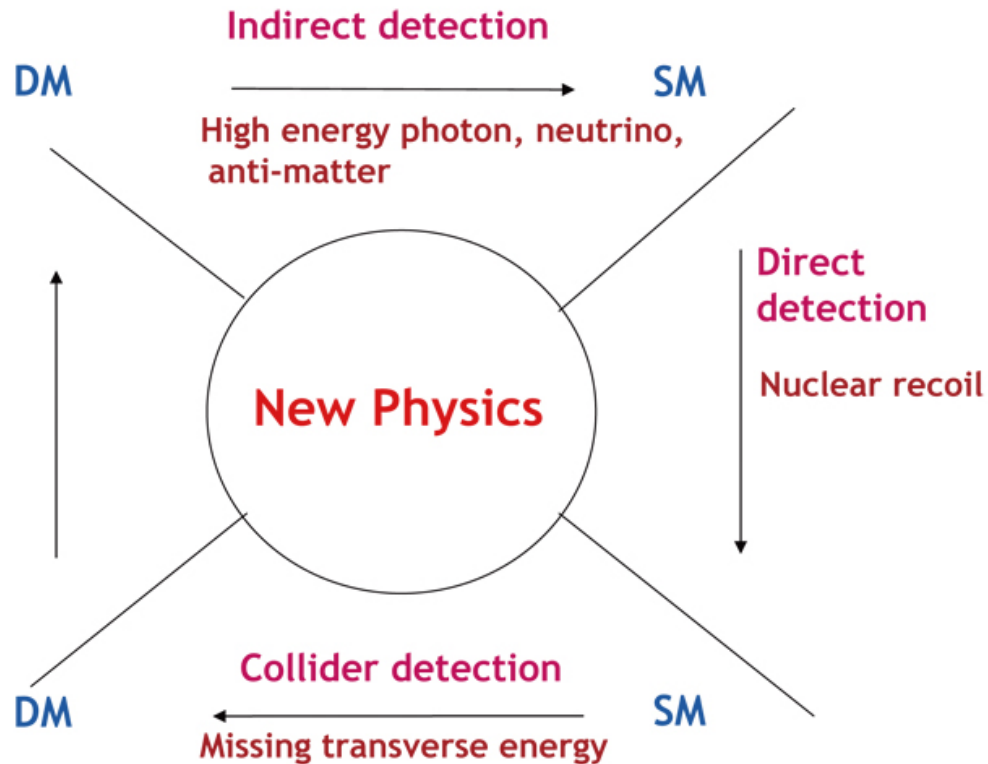
Left: Galactic Center

Right: The Large Magellanic Cloud



Comparison of other methods

- Direct detection:
Signal for DM particle $\ll$ noise,
mostly builded underground
- Collider detection:
High energy input, can be studied
in detail



3 current categories in DM detection

Conclusion

Conclusion

- Occurs when charged velocity $\gg$ light speed in medium
- Can be detected by Cherenkov counter
- Useful in astrophysics (CTA, DM detection etc)

Thanks for listening and Q&A!